

Predictive Engineering and Computational Sciences

# A Streamline-Upwind Petrov-Galerkin Finite Element Scheme for Non-Ionized Hypersonic Flows in Thermochemical Nonequilibrium

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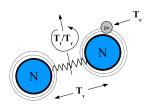


## Acknowledgments

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- 2 This work was supported in part by a cooperative agreement with the Predictive Engineering and Computational Sciences (PECOS) Center at The University of Texas at Austin.

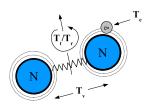
- Background & Motivation
- Physical Modeling
  - Governing Equations
  - Thermochemistry
- Finite Element Formulation
- Results
  - Inviscid Thermal Nonequilibrium Chemically Reacting Flow
  - Viscous Thermal Equilibrium Chemical Reacting Flow
- Near-term Effort

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- Molecules and atoms can store energy in various modes.
- At hypersonic conditions these modes may not be in equilibrium, resulting in thermal nonequilibrium.
- The physical models and governing equations for flows in thermochemical nonequilibrium have been simulated previously with finite difference and finite volume techniques.
- In this work we review the physical models and implement the first known SUPG finite element scheme for hypersonic flows in thermochemical nonequilibrium.

#### **Governing Equations**

 Extension from a single-species calorically perfect gas to a reacting mixture of thermally perfect gases requires species conservation equations and additional energy transport mechanisms.

$$\begin{split} &\frac{\partial \rho}{\partial t} + \boldsymbol{\nabla} \cdot (\rho \ \boldsymbol{u}) = 0 \\ &\frac{\partial \rho \boldsymbol{u}}{\partial t} + \boldsymbol{\nabla} \cdot (\rho \boldsymbol{u} \boldsymbol{u}) = -\boldsymbol{\nabla} P + \boldsymbol{\nabla} \cdot \boldsymbol{\tau} \\ &\frac{\partial \rho E}{\partial t} + \boldsymbol{\nabla} \cdot (\rho H \boldsymbol{u}) = -\boldsymbol{\nabla} \cdot \dot{\boldsymbol{q}} + \boldsymbol{\nabla} \cdot (\boldsymbol{\tau} \boldsymbol{u}) \end{split}$$

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$$\frac{\partial \rho \boldsymbol{u}}{\partial t} + \nabla \cdot (\rho \boldsymbol{u} \boldsymbol{u}) = -\nabla P + \nabla \cdot \boldsymbol{\tau}$$

$$\frac{\partial \rho E}{\partial t} + \nabla \cdot (\rho H \boldsymbol{u}) = -\nabla \cdot \dot{\boldsymbol{q}} + \nabla \cdot (\boldsymbol{\tau} \boldsymbol{u}) + \nabla \cdot \left(\rho \sum_{s=1}^{ns} h_s \mathcal{D}_s \nabla c_s\right)$$

 Problem class may also require a multitemperature thermal nonequilibrium option.

$$\frac{\partial \rho e_V}{\partial t} + \boldsymbol{\nabla} \cdot (\rho e_V \boldsymbol{u}) = -\boldsymbol{\nabla} \cdot \dot{\boldsymbol{q}}_V + \boldsymbol{\nabla} \cdot \left(\rho \sum_{s=1}^{ns} e_{Vs} \mathcal{D}_s \boldsymbol{\nabla} c_s\right) + \dot{\omega}_V$$

## Thermodynamics & Transport Properties

 Thermochemistry models must be extended for a mixture of vibrationally and electronically excited thermally perfect gases.

$$\begin{split} e^{\mathsf{int}} = & e^{\mathsf{trans}} + e^{\mathsf{rot}} + e^{\mathsf{vib}} + e^{\mathsf{elec}} + h^0 \\ = & \sum_{s=1}^{ns} c_s e^{\mathsf{trans}}_s \left( T \right) + \sum_{s=mol} c_s e^{\mathsf{rot}}_s \left( T \right) + \\ & \sum_{s=mol} c_s e^{\mathsf{vib}}_s \left( T_V \right) + \sum_{s=1}^{ns} c_s e^{\mathsf{elec}}_s \left( T_V \right) + \sum_{s=1}^{ns} c_s h^0_s \end{split}$$

Here we have assumed that  $T^{\text{trans}} = T^{\text{rot}} = T$  and  $T^{\text{vib}} = T^{\text{elec}} = T_V$ 

- Additional transport property models are required. In this work we use
  - species viscosity given by Blottner curve fits,
  - species conductivities determined from an Eucken relation,
  - ▶ mixture transport properties computed via Wilke's mixing rule, and
  - ► mass diffusion currently treated by assuming constant Lewis number.

#### **Chemical Kinetics**

We consider r general reactions of the form

$$N_2 + \mathcal{M} \rightleftharpoons 2N + \mathcal{M}$$
...
$$N_2 + O \rightleftharpoons NO + N$$
...

· The reactions are of the form

$$\mathcal{R}_r = k_{br} \prod_{s=1}^{ns} \left( \frac{\rho_s}{M_s} \right)^{\beta_{sr}} - k_{fr} \prod_{s=1}^{ns} \left( \frac{\rho_s}{M_s} \right)^{\alpha_{sr}}$$

where  $\alpha_{sr}$  and  $\beta_{sr}$  are the stoichiometric coefficients for reactants and products

· The source terms are then

$$\dot{\omega}_s = M_s \sum_{r=1}^{nr} (\alpha_{sr} - \beta_{sr}) (\mathcal{R}_{br} - \mathcal{R}_{fr})$$

#### Kinetic Rates

 $\bullet$  The forward rate coefficients are defined with a modified Arrhenius law as a function of some temperature  $\bar{T}$ 

$$k_{fr}\left(\bar{T}\right) = C_{fr}\bar{T}^{\eta_r}\exp\left(-E_{ar}/R\bar{T}\right)$$

where the rate constants are determined empirically.

ullet The corresponding backward rate coefficient can be found using the principle of detailed balance and the equilibrium constant  $K_{eq}$ 

$$K_{eq} = \frac{k_{fr}}{k_{br}}$$

- In thermal equilibrium  $\bar{T}=T.$  We are currently using CANTERA in this regime.
- In thermal nonequilibrium  $\bar{T}=\bar{T}\left(T,T_{V}\right)$  and typical hackery ensues.

#### **Energy Exchange**

$$\dot{\omega}_V = \dot{Q}_v + \dot{Q}_{\mathsf{transfer}}$$

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We adopt the Landau-Teller vibrational energy exchange model

$$\dot{Q}_{s}^{\text{tr-vib}} = \rho_{s} \frac{\hat{e}_{s}^{\text{vib}} - e_{s}^{\text{vib}}}{\tau_{s}^{\text{vib}}} \tag{1}$$

where  $\hat{e}_s^{\rm vib}$  is the species equilibrium vibrational energy and the vibrational relaxation time  $au_s^{\rm vib}$  is given by Millikan and White

$$\tau_s^{\mathsf{vib}} = \frac{\sum_{r=1}^{ns} \chi_r}{\sum_{r=1}^{ns} \chi_r / \tau_{sr}^{\mathsf{vib}}}, \quad \chi_r = c_r \frac{M}{M_r}, \quad M = \left(\sum_{s=1}^{ns} \frac{c_s}{M_s}\right)^{-1}$$

and

$$\begin{split} \tau_{sr}^{\text{vib}} &= \frac{1}{P} \exp \left[ A_{sr} \left( T^{-1/3} - 0.015 \mu_{sr}^{1/4} \right) - 18.42 \right] \\ A_{sr} &= 1.16 \times 10^{-3} \mu_{sr}^{1/2} \theta_{vs}^{4/3}, \;\; \mu_{sr} = \frac{M_s M_r}{M_s + M_r} \end{split}$$

## Vibrational Energy Production and Energy Exchange

$$\dot{\omega}_V = \dot{Q}_v + \dot{Q}_{ ext{transfer}}$$

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When molecular species are created in the gas at rate  $\dot{\omega}_s$ , they contribute vibrational/electronic energy at the rate

$$\dot{Q}_{vs} = \dot{\omega}_s \left( e_s^{\rm vib} + e_s^{\rm elec} \right)$$

so the net vibrational energy production rate is

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 (2)

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$$\dot{Q}_v = \sum_{s=1}^{ns} \dot{\omega}_s \left( e_s^{\mathsf{vib}} + e_s^{\mathsf{elec}} \right) \tag{2}$$

Combining (1) and (2) yields the desired net vibrational energy source term

$$\dot{\omega}_{V} = \sum_{s=1}^{ns} \dot{Q}_{s}^{\text{tr-vib}} + \sum_{s=1}^{ns} \dot{\omega}_{s} \left( e_{s}^{\text{vib}} + e_{s}^{\text{elec}} \right)$$

$$rac{\partial oldsymbol{U}}{\partial t} + rac{\partial oldsymbol{F}_i}{\partial x_i} = rac{\partial oldsymbol{G}_i}{\partial x_i} + \dot{oldsymbol{\mathcal{S}}}$$

$$rac{\partial oldsymbol{U}}{\partial t} + \left(oldsymbol{A}_i^c + oldsymbol{A}_i^P
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ight) + \dot{oldsymbol{\mathcal{S}}}$$

$$\frac{\partial \boldsymbol{U}}{\partial t} + \left(\boldsymbol{A}_{i}^{c} + \boldsymbol{A}_{i}^{P}\right) \frac{\partial \boldsymbol{U}}{\partial x_{i}} = \frac{\partial}{\partial x_{i}} \left(\boldsymbol{K}_{ij} \frac{\partial \boldsymbol{U}}{\partial x_{j}}\right) + \dot{\boldsymbol{S}}$$

Find U satisfying the essential boundary and initial conditions such that

$$\begin{split} \int_{\Omega} \left[ \boldsymbol{W} \cdot \left( \frac{\partial \boldsymbol{U}}{\partial t} + \boldsymbol{A}_{i}^{P} \frac{\partial \boldsymbol{U}}{\partial x_{i}} - \dot{\boldsymbol{\mathcal{S}}} \right) + \frac{\partial \boldsymbol{W}}{\partial x_{i}} \cdot \left( \boldsymbol{K}_{ij} \frac{\partial \boldsymbol{U}}{\partial x_{j}} - \boldsymbol{A}_{i}^{c} \boldsymbol{U} \right) \right] \ d\Omega \\ + \sum_{e=1}^{n_{el}} \int_{\Omega_{e}} \boldsymbol{\tau}_{\mathsf{SUPG}} \frac{\partial \boldsymbol{W}}{\partial x_{k}} \cdot \boldsymbol{A}_{k} \left[ \frac{\partial \boldsymbol{U}}{\partial t} + \boldsymbol{A}_{i} \frac{\partial \boldsymbol{U}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left( \boldsymbol{K}_{ij} \frac{\partial \boldsymbol{U}}{\partial x_{j}} \right) - \dot{\boldsymbol{\mathcal{S}}} \right] \ d\Omega \\ + \sum_{e=1}^{n_{el}} \int_{\Omega_{e}} \nu \left( \frac{\partial \boldsymbol{W}}{\partial x_{i}} \cdot g^{ij} \frac{\partial \boldsymbol{U}}{\partial x_{j}} \right) \ d\Omega - \oint_{\Gamma} \boldsymbol{W} \cdot (\boldsymbol{g} - \boldsymbol{f}) \ d\Gamma = 0 \end{split}$$

for all W in an appropriate function space.

#### Stabilization Parameters

$$\nu = \left[ \frac{\left\| \frac{\partial \boldsymbol{U}}{\partial t} + \boldsymbol{A}_i \frac{\partial \boldsymbol{U}}{\partial x_i} - \frac{\partial}{\partial x_i} \left( \boldsymbol{K}_{ij} \frac{\partial \boldsymbol{U}}{\partial x_j} \right) \right\|_{\boldsymbol{A}_0^{-1}}^2}{\left( \Delta \boldsymbol{U}_h \right)^T \boldsymbol{A}_0^{-1} \Delta \boldsymbol{U}_h + g^{ij} \left( \frac{\partial \boldsymbol{U}_h}{\partial x_i} \right)^T \boldsymbol{A}_0^{-1} \frac{\partial \boldsymbol{U}_h}{\partial x_j}} \right]^{1/2}$$

 $\boldsymbol{\tau}_{\text{SUPG}} = \text{diag}\left(\tau_{c,s}, \tau_{m,i}, \tau_{E}, \tau_{e_{V}}\right)$ 

where  $\tau_c$ ,  $\tau_{m,j}$ ,  $\tau_E$ , and  $\tau_{e_V}$  are given by

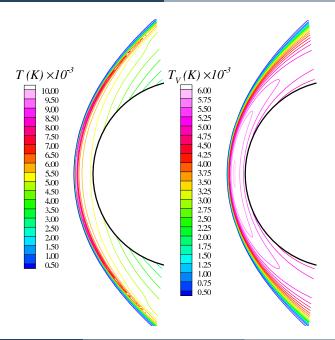
$$\begin{split} \tau_{c,s} &= \left[ \left( \frac{2 \left( \| \boldsymbol{u} \| + c \right)}{h_{\mathrm{SUPG}}} \right)^2 + \left( \frac{4 \mathcal{D}_s}{h_{\mathrm{SUPG}}^2} \right)^2 + \nu^2 \right]^{-1/2} \\ \tau_{m,j} &= \left[ \left( \frac{2 \left( \| \boldsymbol{u} \| + c \right)}{h_{\mathrm{SUPG}}} \right)^2 + \left( \frac{4 \mu}{\rho h_{\mathrm{SUPG}}^2} \right)^2 + \nu^2 \right]^{-1/2} \\ \tau_E &= \left[ \left( \frac{2 \left( \| \boldsymbol{u} \| + c \right)}{h_{\mathrm{SUPG}}} \right)^2 + \left( \frac{4 k}{\rho c_p h_{\mathrm{SUPG}}^2} \right)^2 + \nu^2 \right]^{-1/2} \\ \tau_{ev} &= \left[ \left( \frac{2 \left( \| \boldsymbol{u} \| + c \right)}{h_{\mathrm{SUPG}}} \right)^2 + \left( \frac{4 k_v}{\rho C_v^{\mathrm{vib}} h_{\mathrm{SUPG}}^2} \right)^2 + \nu^2 \right]^{-1/2} \end{split}$$

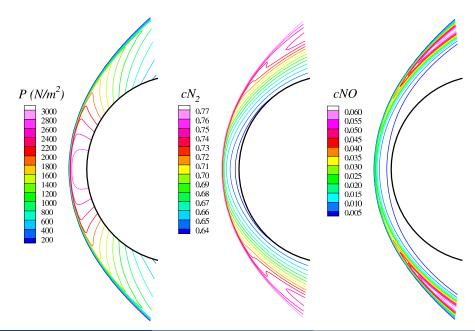
#### Inviscid Cylinder

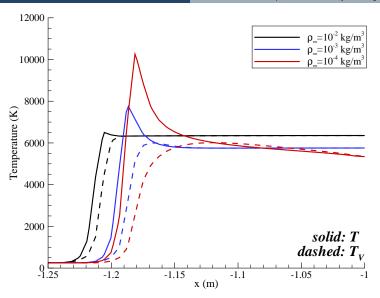
- Inviscid flow in thermochemical nonequilibrium
- 5 species air (N<sub>2</sub>, O<sub>2</sub>, NO, N, O)
- 5 reaction model with Park 1990 rates

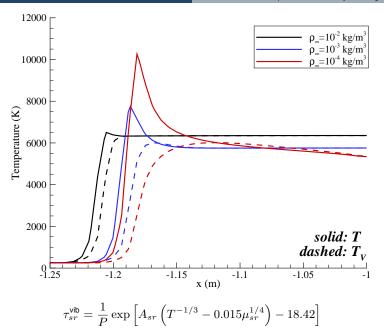
$$cN_{2,\infty} = 0.78, cO_{2,\infty} = 0.22$$
  
 $U_{\infty} = 5,500 \,\mathrm{m/sec}$   
 $\rho_{\infty} = 10^{-2} - 10^{-4} \,\mathrm{kg/m^3}$   
 $T_{\infty} = 250 \,\mathrm{K} = T_{V,\infty}$ 

 Landau-Teller vibrational energy relaxation model, with Millikan and White vibrational relaxation time







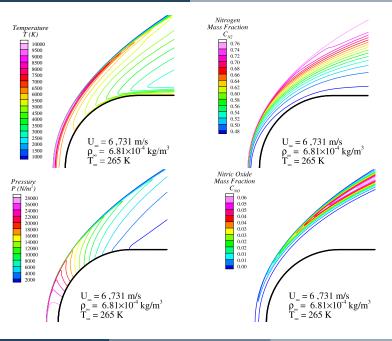


## 2D Extended Cylinder

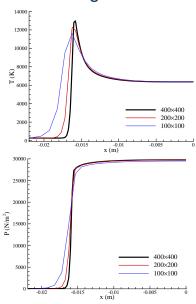
- Laminar flow in thermal equilibrium
- No-slip, adiabatic, noncatalytic wall
- Chemical nonequilibrium, 5 species air (N<sub>2</sub>, O<sub>2</sub>, NO, N, O)
- 5 reaction model with Park 1990 rates

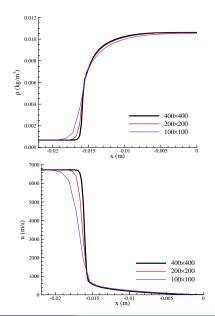
$$c{\sf N}_{2,\infty} = 0.78, c{\sf O}_{2,\infty} = 0.22$$
  $U_{\infty} = 6,731\,{\rm m/sec}$   $ho_{\infty} = 6.81 \times 10^{-4}\,{\rm kg/m^3}$   $T_{\infty} = 265\,{\rm K}$ 

- Blottner/Wilke/Eucken with constant Lewis number Le=1.4 for transport properties
- · Mesh, iterative convergence
- FIN-S/DPLR comparison
- Weak & Strong Scaling

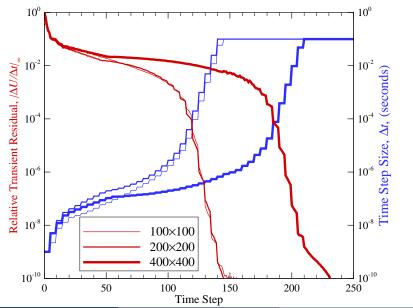


# Mesh Convergence

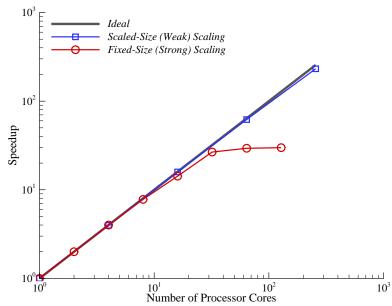




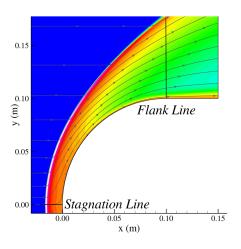
## Iterative Convergence

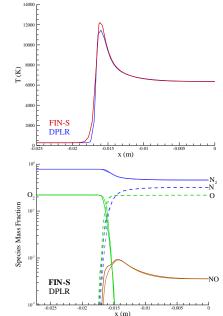


# Speedup

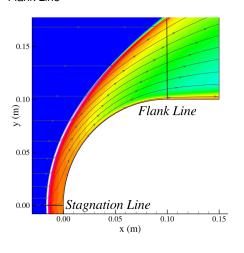


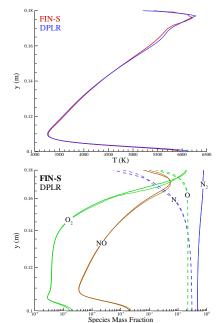
# Code-to-Code Comparison – Stagnation Line





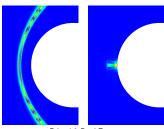
# Code-to-Code Comparison – Flank Line



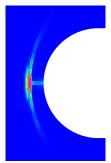


#### Additional Focus Areas

- Physics Modeling
  - Weakly Ionized Flows
  - ► Surface Catalycity
  - ► Additional Boundary Conditions
- 2 Ablation coupling
- 3 Adjoints
  - Sensitivity analysis
  - Adaptivity



Primal & Dual Error



Combined Qol Error

Thank you!

Questions?